# Observational Effects of Strong Gravity in Vicinity of Supermassive Black Holes

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### Summary of the talk:

- Supermassive black holes as driving engines of Active
  Galactic Nuclei (AGN) and Quasars
- □ Standard model of accretion disk
- X-ray radiation from accretion disk around a supermassive black hole (BH): observational effects
- Our recent investigations and results
- □ Conclusions

## Active Galaxies and Quasars

 small highly variable and very bright core embedded in an otherwise typical galaxy

		· .
	Normal Spiral G	alaxy
_	*	
Galaxy v	with an Active Galac	tic Nucleus (AGN)
	*	* ·
	Quasar	

#### features:

- □ 10% of all galaxies
- 10<sup>4</sup> times higher luminosity than typical galaxies
- $\Box$  tiny volumes (« 1 pc<sup>3</sup>)
- radiation in broad range: from γ-rays to radio waves
- very small angular size depending on wavelength
- strong and sometimes very broad emission lines
- variability
- polarization
- □ radio emission

# Unified model of AGN

### **Classification of AGN**

- □ radio-loud galaxies
- □ radio-quiet galaxies
- □ broad line radio galaxies
- □ narrow line radio galaxies
- □ optically violently variable quasars
- □ BL Lac objects (blazars)
- □ Seyfert I galaxies
- □ Seyfert II galaxies
- LINERs (Low-Ionization Nuclear Emission Regions)

### Urry & Padovani 1995, PASP, Vol. 107, No. 715, 803.











PHYSICS

**KINEMATICS** 

(Martin, E. 2000, ApJ, 545, 63)

### Some examples of AGN







### Space-time geometry in vicinity of BH

Schwarzschild metric:

$$ds^{2} = dt^{2} \left(1 - \frac{1}{r}\right) - dr^{2} \left(1 - \frac{1}{r}\right)^{-1} - r^{2} \left(d\theta^{2} + \sin^{2}\theta d\phi^{2}\right)$$

Schwarzschild radius:

$$r_s = \frac{2GM}{c^2}$$

Kerr metric:

$$ds^{2} = -\left(1 - \frac{2Mr}{\Sigma}\right) dt^{2} - \frac{4Mar}{\Sigma} \sin^{2}\theta \, dt \, d\phi + \frac{A}{\Sigma} \sin^{2}\theta \, d\phi^{2} + \frac{\Sigma}{\Delta} \, dr^{2} + \Sigma \, d\theta^{2}$$

where

$$\Sigma = r^2 + a^2 \cos^2 \theta$$
,  $\Delta = r^2 + a^2 - 2Mr$ ,  $A = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta$ .

#### Standard model of accretion disk

- Optically thick and geometrically thin disk
- □ power law of disk emissivity:

 $I_P(X,Y) = \varepsilon_0 \cdot r^q(X,Y) \cdot g^4(X,Y)$ 

□ black-body emissivity law:

$$I_P(X,Y;E) = \frac{2E^3}{h^2c^2} \frac{1}{e^{E/kT(X,Y)} - 1}$$

□ modifiedblack-body emissivity law:

$$I_P(X,Y;E) \propto x^3 e^{-x}, \quad x = rac{E}{kT(X,Y)}$$

□ effective temperature of radiation:  $10^7 - 10^8$  K

Radial distribution of surface temperature (Shakura & Sunyaev 1973, *Astron. Astrophys*, **24**, 337)



The distribution of the temperature as a function of the radius *R* along the direction of the disk rotation, given for two different values of angular momentum *a*. Negative values of *R* correspond to the approaching and positive values to the receding side of the disk.



#### X-ray radiation from accretion disks of AGN

- **in continnum**: 0.1 100 keV
- □ soft and hard component
- variations: from several part of an hour until several days

#### **<u>in Fe Kα line:</u>**

- □ broad emission line on 6.4 keV
- asymetric profile with narrow bright blue peak and wide faint red peak
- Line width corresponds to velocity:
  - $v \sim 80000 100000 \text{ km/s} (MCG-6-30-15)$
  - $v \sim 48000 \text{ km/s} (\text{MCG-5-23-16})$
  - $v \sim 20000 30000$  km/s (many other AGN)
- □ variability of both: line shape and intensity



MCG-6-30-15, ASCA SIS (Tanaka et al. *Nature*, 1995, **375**, 659.)

### Ray-tracing in Kerr metric

$$\begin{split} &\pm \int_{r_{\rm em}}^{\infty} \frac{dr}{\sqrt{R(r,\lambda,q)}} = \pm \int_{\theta_{\rm em}}^{\theta_{\rm obs}} \frac{d\theta}{\sqrt{\Theta(\theta,\lambda,q)}} \\ &R(r,\lambda,q) = \left\{ (r^2 + a^2 - a\lambda)^2 - \Delta[(\lambda - a)^2 + q^2] \right\} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right] \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right] \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right] \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right] \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (\lambda\cot\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 \right]^{1/2} \\ &\Theta(\theta,\lambda,q) = \left[ q^2 + (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)^2 - (a\cos\theta)$$

# Our recent investigations and results

- Modeling of emission of accretion disk around supermassive BH using numerical simulations based on a ray-tracing method
- Investigation of the effects of strong gravitational field by comparison between modeled and observed iron Kα line profiles
- Scanning the innermost parts of accretion disks by gravitational microlensing: variations of the Kα line profile

Numerical simulations of a highly inclined accretion disk ( $i=75^{\circ}$ ) for different values of angular momentum parameter *a* (left) and the corresponding profiles of the Fe Ka line (right)



Numerical simulations of an accretion disk in Schwarzschild metric for different inclination angles i (left) and the corresponding profiles of the Fe K $\alpha$  line (right)



Numerical simulations of an accretion disk in Kerr metric with angular momentum parameter a = 0.998 for different inclination angles *i* (left) and the corresponding profiles of the Fe Ka line (right)





The Fe K $\alpha$  line profile from Seyfert I galaxy MCG-6-30-15 observed by the ASCA satellite (Tanaka, Y. et al, 1995, *Nature*, **375**, 659). The solid line shows the modeled profile expected from an accretion disk extending between 6 and 20  $R_g$  around Schwarzschild BH.



#### Gravitational microlensing influence on disk emission

amplification of a point-like microlens:  $A(X, Y) = \frac{u^2(X, Y) + 2}{u(X, Y)\sqrt{u^2(X, Y) + 4}}$ 

amplification of a caustic microlens:  $A(X, Y) = A_0 + \frac{K}{\sqrt{\kappa(\xi - \xi_c)}} \cdot H(\kappa(\xi - \xi_c))$ 

amplified brightness:  $I_p = \varepsilon(r)g^4(X, Y)\delta(x - g(X, Y))A(X, Y), \quad g = v_{obs}/v_0$ 

total observed flux: 
$$F(x) = \int_{\text{image}} I_p(x) d\Omega$$

where  $d\Omega$  is the solid angle subtended by the disc in the observer's sky





# Conclusions

- Comparisons between the modeled and observed iron Kα line profiles allow us to determine the parameters of the line emitting region.
- Two of them are of especial importance for investigating the strong gravitational field of AGN: mass of central BH and its angular momentum.
- Our results show that these parameters have significant influence on Fe Kα line profile and thus, allow us to determine the space-time geometry (metric) in vicinity of the central BH of AGN.
- Other parameters can give us information about the plasma conditions in these regions.
- □ Gravitational microlensing is a useful tool for revealing the structure and physics of the innermost parts of accretion disks in AGN, close to their central supermassive black holes.

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## Thank you for attention! Lpsuk you for attention!