

# SUPERSYMMETRY BREAKING

BORUT BAJC

JSI, LJUBLJANA

BW 2007

WORK DONE WITH  
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(ALEJANDRA MELFO)

# OUTLINE

① INTRODUCTION

② THE PHILOSOPHY

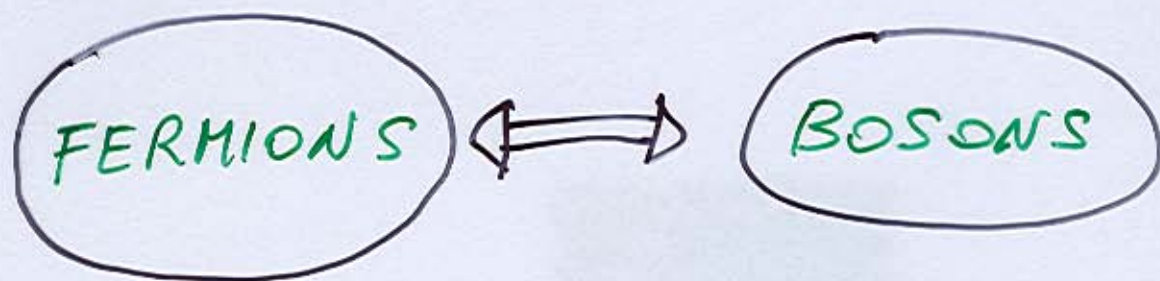
③ AN EXPLICIT EXAMPLE

- SUSY BREAKING

- MEDIATION

④ CONCLUSIONS

# SUPERSYMMETRY:



$$e \rightarrow E = \tilde{e} + \theta e + \theta\theta F$$

Labels for the equation above:

- $e$ : FERMION  $S = 1/2$
- $E$ : SUPERFIELD
- $\tilde{e}$ : BOSON  $S = 0$
- $\theta e$ : FERMION  $S = 1/2$
- $\theta\theta F$ : AUXILIARY FIELD

$\theta$  ... GRASSMAN VARIABLE

$$F \neq 0 \iff \text{SUSY}$$

①

SUPERSYMMETRY (SUSY)

MOTIVATED BY (~~UNIFICATION~~)

- GAUGE COUPLING UNIFICATION
- HIERARCHY PROBLEM
- DYNAMICAL ELECTROWEAK SYMMETRY BREAKING
- DARK MATTER

SUSY PART OF MSSM (W) KNOWN

~~SUSY~~ PART OF MSSM UNKNOWN

PHENOMENOLOGICAL CONSTRAINTS :

DIRECT SEARCHES  $\Rightarrow \tilde{m} \gtrsim 100 \text{ GeV}$

FCNC  $\Rightarrow$  FLAVOUR STRUCTURE OF  
SFERMIONS SIMILAR TO  
FERMIONS ( $V_{CKM}$ )

① CAN WE BREAK SUSY  
IN MSSM?

$$F_{H, \bar{H}} \neq 0 \quad (\text{THE ONLY POSSIBILITY})$$

$$\int d^2\theta y_e H \cdot \cancel{L} e^c \rightarrow y_e F_H \tilde{e} \tilde{e}^c$$

$$V = \sum_i \left| \frac{\partial W}{\partial \varphi_i} \right|^2 = y_e^2 |H|^2 (|\tilde{e}|^2 + |\tilde{e}^c|^2) + \dots$$

$$\begin{pmatrix} \tilde{e}^* & \tilde{e}^c \end{pmatrix} \begin{pmatrix} y_e^2 v^2 & y_e^2 F \\ y_e^2 F & y_e^2 v^2 \end{pmatrix} \begin{pmatrix} \tilde{e} \\ \tilde{e}^{c*} \end{pmatrix}$$

$$m_{1,2}^2 = y_e^2 (v^2 \pm F)$$

$\Downarrow$

$$m_1^2 + m_2^2 = 2 m_e^2$$

IMPOSSIBLE!

SUCH ~~OPERATORS~~  
MASS SUM RULE

ALWAYS IN SUSY  
AT TREE ORDER.

WAYS OUT:

a) USE OF LOCAL SUSY (SUGRA)

(MEDIATION)

b) SUSY BREAKING BY LOOPS

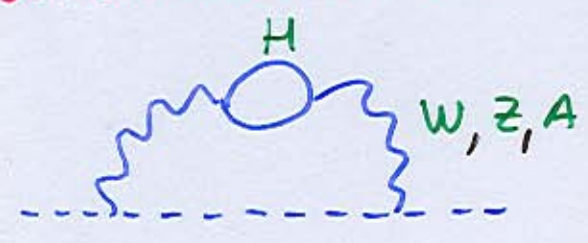
IN BOTH CASES HIGHER DIMENSIONAL  
~~NONRENORMALIZABLE~~ OPERATORS

(BUT IN b) CALCULABLE)

② CAN ONE ADD JUST ONE EXTRA FIELD WITH  $F \neq 0$ ? **No!**

$\int d^2\theta H\bar{H}(\mu+X) \Rightarrow m_{\tilde{H}} = \mu$  (HIGGSINO)  
 $m_{H_1}^2 = \mu^2 - F$  (SM HIGGS)  
 $m_{H_2}^2 = \mu^2 + F$

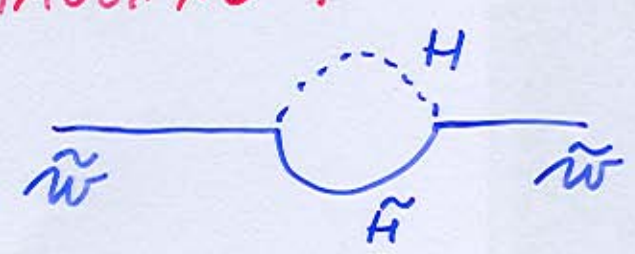
OTHER FERMIONS:



$\tilde{m}^2 \sim \left(\frac{\alpha}{4\pi}\right)^2 \left(\frac{F}{\mu}\right)^2 \ll F$

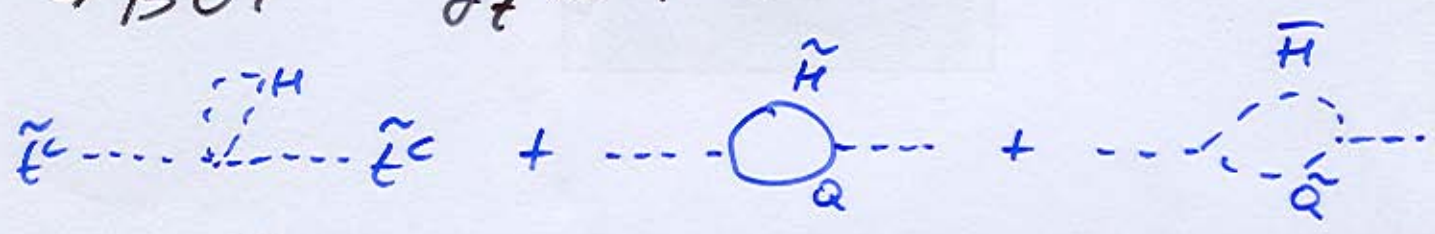
$F \sim \mu^2$  LARGE ( $m_{H_1}^2 \sim m_w^2$ )

GAUGINO:



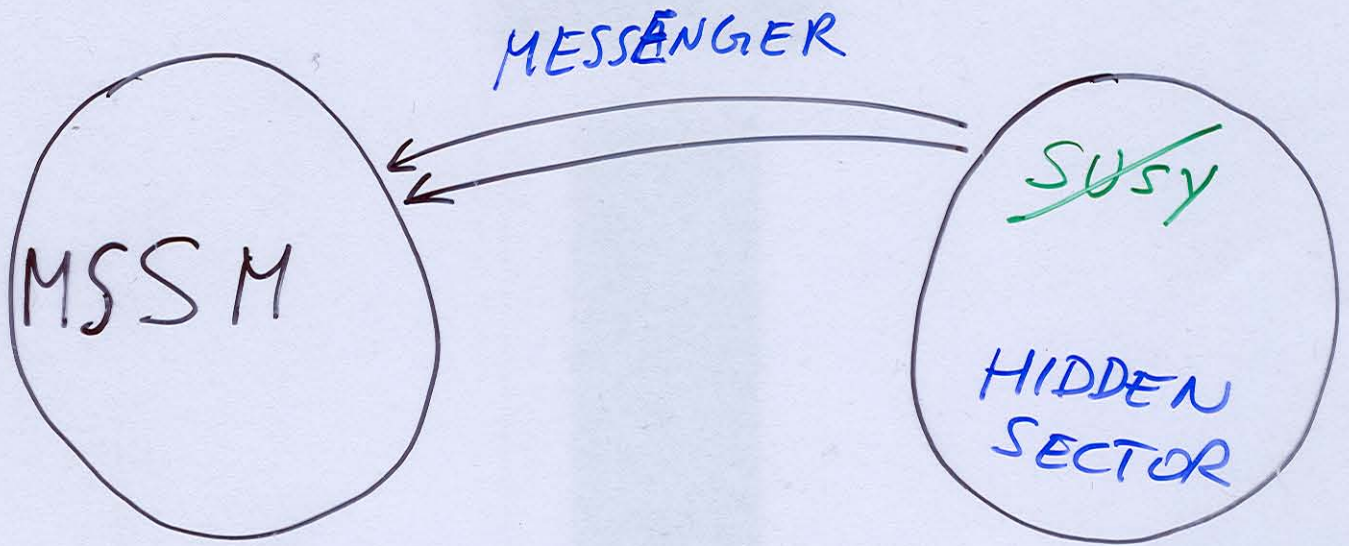
→ BUT NO 1-LOOP MASS FOR GLUINO!  
**1<sup>ST</sup> PROBLEM**

→ BUT  $y_t \sim 1$ :



$\Rightarrow m_{\tilde{t}^c}^2 \approx -y_t^2 F$

LARGE AND NEGATIVE!  
**2<sup>ND</sup> PROBLEM**





### ③ FEW SHORT COMMENTS ON KNOWN MODELS OF SUSY :

#### (A) GRAVITY MEDIATION

MESSANGER  $\Rightarrow$  GRAVITY

- GRAVITY ALREADY THERE
- JUST A PARAMETRIZATION,  
SOFT TERMS DEPEND ON ~~THESE~~  
ASSUMPTIONS ON KÄHLER

#### CMSSM :

EVERYTHING CALCULABLE  
IN TERMS OF FEW PARAMETERS

$(m_0, m_{1/2}, A, B, \dots)$ ,

BUT THEORETICALLY NOT MOTIVATED!

WHY SHOULD KÄHLER BE CANONICAL?

# (B) ANOMALY MEDIATION

AGAIN CONNECTED WITH  
GRAVITY, ALWAYS THERE

VERY PREDICTIVE  $\Rightarrow$  WRONG

~~THESE ARE THE PROBLEMS~~

$$m^2 \sim \beta_g \quad (\text{RG INVARIANT})$$

$\Rightarrow$  SLEPTONS GET  $m^2 < 0$

NEED TO COMBINE WITH  
OTHER SOURCES OF SUSY:

BACK TO OLD UNCERTAINTIES

# (C) GAUGE MEDIATION

ALREADY SEEN AN UNSUCCESSFUL

EXAMPLE WITH  $-F_X \neq 0$

- X COUPLED TO,  $H \bar{H}$

$$F_X \neq 0$$

EXTRA  $q, \bar{q}, \ell, \bar{\ell}$   
GETS ~~MSSM~~ MASSES

$\Rightarrow$  MSSM  
SM  
GAUGE  
BOSONS

$$W = \lambda_q X \bar{q} q + \lambda_\ell X \bar{\ell} \ell + W_{\text{MSSM}}$$

APPARENTLY FLAVOUR CONSERVING  
(GAUGE INTERACTIONS FLAVOUR BLIND)

BUT TO SAVE GAUGE COUPLING UNIFICATION

$$5 \sim (q, \ell) \quad \bar{5} \sim (\bar{q}, \bar{\ell})$$

" $d^c$ "      " $L$ "

$\Rightarrow \tilde{Y} H \ell e^c$  POSSIBLE  
(EXTRA FLAVOUR MIXING!)

$$\Rightarrow \tilde{Y} \ll 1$$

NOT ~~MUCH~~ BETTER THAN SUGRA!

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EXTRA  $q, \bar{q}, l, \bar{l}$   
GETS MSSM MASSES

$\Rightarrow$  SM GAUGE BOSONS  
MSSM

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④ SUSY NATURALLY TIED  
WITH GUTS  
(GAUGE COUPLING UNIFICATION)

CAN WE GET ANY  
INSIGHT IF WE CONNECT  
THE FIELD  $X$  THAT  
BREAKS SUSY ( $F_x \neq 0$ )  
WITH A GUT MULTIPLY ?

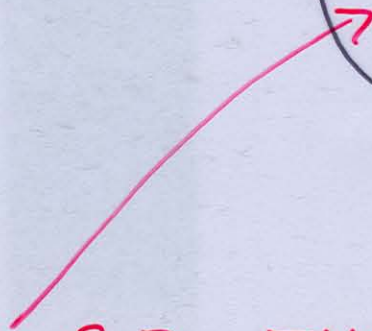
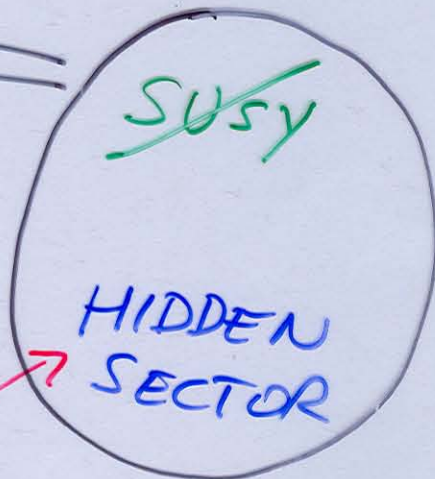
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CAN WE GET ANY  
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THE FIELD  $X$  THAT  
BREAKS SUSY ( $F_x \neq 0$ )  
WITH A GUT MULTIPLY ?

MAINLY  
GRAVITY



MESSENGER



CAN IT BE THE HIGGS  
SECTOR IN A GUT?

(FOR EX.  $24_H$  IN  $SU(5)$ )

# RULES OF THE GAME :

① NO EXTRA INTERACTIONS  
EXCEPT GUT

(NO ROOM FOR DYNAMICAL  
SUPERSYMMETRY BREAKING  
OR FAYET-ILIOPOULOS)

AND GRAVITY

(= MESSENGER)

② NO SINGLET

(NO ROOM FOR O'RAIFEARTAIGH)



CAN THE SAME FIELD  
BREAK BOTH  
GUT AND SUSY ?

TYPICALLY CONSIDERED  
DIFFICULT,  
DIFFERENT ARGUMENTS  
AGAINST IT

# PROBLEM

LOCAL SUSY (SUGRA)

=

GLOBAL SUSY +  $\frac{1}{M_{Pl}}$  CORRECTIONS  
(TYPICALLY SMALL)



~~SUSY~~ IN LOCAL SUSY

JUST GENERALIZATION OF

~~SUSY~~ IN GLOBAL SUSY

(O'RAIFEARTAIGH OR FAYET-ILIPOULOS)



# PROBLEM

~~SUSY~~



GOLDSTINO MASSLESS



PARTNERS IN GUT & SUSY  
MULTIPLY EASILY LIGHT



SPOILS GAUGE COUPLING  
UNIFICATION

$$\Sigma = \sigma \begin{pmatrix} 2 & & & \emptyset \\ & 2 & & \\ & & 2 & \\ \emptyset & & & -3 \\ & & & & -3 \end{pmatrix} + \begin{pmatrix} \Sigma_8 & & & \\ & & & \\ & & & \\ & & & \\ & & & \Sigma_3 \end{pmatrix}$$

SIMILAR TO X BEFORE

$$F_\sigma \neq 0$$

$$\langle \sigma \rangle = M_{\text{GUT}}$$

$\Sigma_8$  ... COLOR OCTET

USUALLY

$\Sigma_3$  ... WEAK TRIPLET

$$M_{3,8} \sim M_{\text{GUT}}$$

WE WILL SEE THAT  
MINIMALITY IMPLIES

$$F_\sigma \neq 0 \iff M_{3,8} \ll M_{\text{GUT}}$$

- ONE FIELD  $\sigma$
- KÄHLER CANONICAL

$$V = e^{\frac{|\sigma|^2}{M_{pe}^2}} \left( \left| W' + \sigma^* \frac{W}{M_{pe}^2} \right|^2 - 3 \frac{|W|^2}{M_{pe}^2} \right)$$

CONSTRAINTS ON  $W(\sigma)$

$$\langle W \rangle = \mu_{3/2} M_{pe}^2$$

$$\langle V \rangle = 0 \iff \langle W' \rangle = \sqrt{3} \mu_{3/2} M_{pe} \left( 1 + \sigma \left( \frac{\langle \sigma \rangle}{M_{pe}} \right) \right)$$

$$\left\langle \frac{\partial V}{\partial \sigma} \right\rangle = 0 \iff \langle W'' \rangle = 2 \mu_{3/2} \left( 1 + \sigma \left( \frac{\langle \sigma \rangle}{M_{pe}} \right) \right)$$

MINIMUM

(NOT MAXIMUM)

$$\iff \left| \langle W''' \rangle \right| \leq \frac{2 \mu_{3/2}}{\sqrt{3} M_{pe}} \left( 1 + \sigma \left( \frac{\langle \sigma \rangle}{M_{pe}} \right) \right)$$

AT  $\langle \sigma \rangle = M_{GUT}$

WE NEED  $W(\sigma)$  WITH THESE

4 CONSTRAINTS  $\implies$  4 FREE PARAMETERS

$$W = m\sigma^2 + \lambda\sigma^3 + \frac{c_1}{M_{pe}}\sigma^4 + \frac{c_2}{M_{pe}^2}\sigma^5$$



PARAMETRIZE

$$W = \frac{1}{M_{pe}^2} \left( a_1 M_{GUT}^3 \sigma^2 + a_2 M_{GUT}^2 \sigma^3 + a_3 M_{GUT} \sigma^4 + a_4 \sigma^5 \right)$$

CLOSE TO  $\sigma = M_{GUT}$  ALL

4 TERMS COMPARABLE

OVRUT, RABY, 83

INTERESTED IN LOW-ENERGY SUSY

$$M_{3/2} \approx \mathcal{O}(\text{TeV}) \ll M_{\text{GUT}}$$

TYPICAL

$$\frac{d^m W}{d\tau^m} \approx \underbrace{a}_{\text{blue scribble}} M_{\text{GUT}}^{3-m} \approx M_{3/2} M_{\text{Pl}}^{2-m}$$

$$m = 0, 1, 2, 3$$

IT FOLLOWS:

① AT LEAST 4 COEFFICIENTS  $a$ 's OF COMPARABLE ORDER

② ~~scribble~~  $a \ll 1$

DETERMINE

$a_1, a_2, a_3, a_4$

TO GET THE MINIMUM OF  $V$

AT  $\sigma = M_{\text{GUT}}$  WITH

$$\langle V \rangle = 0$$

$$\langle W \rangle = m_{3/2} M_{\text{Pl}}^2$$



$$\begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ a_4 \end{pmatrix} = M_{\text{Pl}}^2 \begin{pmatrix} 10 & -6 & \frac{3}{2} & -\frac{1}{6} \\ -20 & 14 & -4 & \frac{1}{2} \\ 15 & -11 & \frac{7}{2} & -\frac{1}{2} \\ -4 & 3 & -1 & \frac{1}{6} \end{pmatrix} \begin{pmatrix} \langle W \rangle / M_{\text{GUT}}^5 \\ \langle W' \rangle / M_{\text{GUT}}^4 \\ \langle W'' \rangle / M_{\text{GUT}}^3 \\ \langle W''' \rangle / M_{\text{GUT}}^2 \end{pmatrix}$$

$$a_i \approx \frac{m_{3/2} M_{\text{Pl}}^4}{M_{\text{GUT}}^5} \ll 1$$



SM SINGLET  $\sigma$  CAN COME

FROM  $\Sigma = 24_H$  IN  $SU(5)$

$$W = m \text{Tr} \Sigma^2 + \lambda \text{Tr} \Sigma^3 +$$
$$+ \frac{b_1^{(1)}}{M_{pe}} \text{Tr} \Sigma^4 + \frac{b_1^{(2)}}{M_{pe}} (\text{Tr} \Sigma^2)^2 +$$
$$+ \frac{b_2^{(1)}}{M_{pe}^2} \text{Tr} \Sigma^5 + \frac{b_2^{(2)}}{M_{pe}^2} (\text{Tr} \Sigma^3)(\text{Tr} \Sigma^2)$$

$$\Sigma = \frac{\sigma}{\sqrt{30}} \begin{pmatrix} 2 & & & & \\ & 2 & & & \\ & & -3 & & \\ & & & -3 & \\ & & & & \dots \end{pmatrix} + \begin{pmatrix} \Sigma_8 & & & & \\ & \ddots & & & \\ & & \ddots & & \\ & & & \ddots & \\ & & & & \Sigma_3 \end{pmatrix}$$

$$a b_i = f_i(m, \lambda, a_\alpha^{(\beta)})$$

↑  
4 PARAMETERS  
(DETERMINED!)

SMALL

↑  
6 PARAMETERS  
(2 LEFT)  
CAN BE  $O(1)$

## MORE PRECISELY

$$a_1 \frac{M_{\text{GUT}}^3}{M_{\text{Pl}}^2} = m$$

$$a_2 \frac{M_{\text{GUT}}^2}{M_{\text{Pl}}^2} = -\frac{\lambda}{\sqrt{30}}$$

$$a_i \ll 1$$

$$a_3 \frac{M_{\text{GUT}}}{M_{\text{Pl}}^2} = \frac{1}{M_{\text{Pl}}} \left( \frac{7}{30} b_1^{(1)} + b_2^{(2)} \right)$$

$$a_4 \frac{1}{M_{\text{Pl}}^2} = \frac{1}{M_{\text{Pl}}^2} \left( -\frac{13}{30\sqrt{30}} b_2^{(1)} - \frac{1}{\sqrt{30}} b_2^{(2)} \right)$$



$$\rightarrow m \ll M_{\text{GUT}}$$

$$\rightarrow \lambda \ll 1$$

$\rightarrow b_i^{(\alpha)}$  COULD BE EVEN  $\mathcal{O}(1)$

PROVIDED THE ABOVE COMBINATIONS  $\ll 1$

## BOTTOM LINE : INTERMEDIATE STATES

$$m_3 \approx \frac{8}{3} b_1^{(1)} \frac{M_{\text{GUT}}^3}{M_{\text{pe}}^2} \rightarrow -\frac{28}{3\sqrt{30}} b_2^{(1)} \frac{M_{\text{GUT}}^3}{M_{\text{pe}}^2}$$

$$m_8 \approx \frac{2}{3} b_1^{(1)} \frac{M_{\text{GUT}}^3}{M_{\text{pe}}^2} + \frac{1}{\sqrt{30}} b_2^{(1)} \frac{M_{\text{GUT}}^3}{M_{\text{pe}}^2}$$

RGE :

$$M_{\text{GUT}}^0 = 10^{16} \text{ GeV}$$

$$M_{\text{GUT}} = M_{\text{GUT}}^0 \left( \frac{M_{\text{GUT}}^0}{(m_3 m_8)^{1/2}} \right)^{1/2}$$

$M_{\text{GUT}}$  INCREASES  $(\approx 10^{17} - 10^{18} \text{ GeV})$

$d=6$  PROTON DECAY SUPPRESSED

$$\tau_p(d=6) = \tau_p^0(d=6) \left( \frac{M_{\text{GUT}}}{M_{\text{GUT}}^0} \right)^4$$

RGE:

$$S_H = \begin{pmatrix} T \\ D \end{pmatrix}$$

$$m_T = m_T^0 \left( \frac{m_3}{m_g} \right)^{5/2}$$

EASILY MAKE  $m_T$  INCREASE

$d=5$  PROTON DECAY EASILY SUPPRESSED

FOR EXAMPLE:

$$b_2^{(i)} \approx 0 \Rightarrow m_3 = 4 m_g$$

$$\tau_p(d=5) = \tau_p^0(d=5) \left( \frac{m_T}{m_T^0} \right)^2$$

$$\Downarrow \\ m_T = 32 m_T^0$$

$$\Downarrow \\ \tau_p \approx 10^3 \tau_p^0$$

PRICE TO PAY: FINE-TUNING

$$m_\sigma \approx m_{3/2}$$

$$m_{3,8} \gg m_{3/2}$$

UNAVOIDABLE  
WITH ONLY  
24<sub>H</sub>

UNLESS FINE-TUNING:

$$W = u_{\sigma} \sigma^2 + m_3 \sigma_3^2 + u_8 \sigma_8^2 + \dots$$

$$m_{\sigma} \ll m_{3,8}$$

~~TRIPLET + OCTET SPLITTING~~

(TRIPLET + OCTET) VS. SINGLET SPLITTING

UNLESS TERMS IN  $W$   
VERY SMALL



$\frac{1}{M_{pe}}$  CORRECTIONS RELEVANT!

~~2021~~ MEDICATION OF 2021

①

GRAVITY

②

X, Y GAUGE BOBONS

③

$\rho^3, \rho^8$

④

LIGHT MIPSES

# ① GRAVITY

$$m_{\tilde{f}} \approx \frac{F}{M_{\text{pl}}} \approx m_{3/2}$$

GRAUGINO:

$$\delta \mathcal{L} \sim \int d^2\theta \frac{c}{M_{\text{pl}}} \text{Tr} (\Sigma W^\alpha W_\alpha) + \text{h.c.}$$

$$\Rightarrow \underbrace{\frac{c}{M_{\text{pl}}} F \text{Tr} (\gamma \lambda \lambda)}_{\text{gaugino mass}} + \underbrace{\frac{c}{M_{\text{pl}}} M_{\text{loop}} \text{Tr} (\gamma F^{\mu\nu} F_{\mu\nu})}_{\Delta \left( \frac{1}{g^2} \right)}$$

$c \ll 1 \Rightarrow$  SAME RESULTS APPLY

$$m_\lambda \ll m_{\tilde{f}}$$

WHAT IF  $c \approx \mathcal{O}(1)$  ?



## RGE CHANGED

$$M_{\text{GUT}} = M_{\text{GUT}}^0 \left( \frac{M_{\text{GUT}}^0}{(\mu_3 \mu_8)^{1/2}} \right)^{1/2} e^{-\left[ \frac{5}{24} \frac{(4\pi)^2}{130} \frac{M_{\text{GUT}}}{M_{\text{Pl}} c} \right]}$$

$\sigma(1)$  FOR  $c=1$

$M_{\text{GUT}}$  DOES NOT CHANGE MUCH

$$M_T = \mu_T^0 \left( \frac{\mu_3}{\mu_8} \right)^{5/2} e^{+26 \left[ \frac{5}{24} \frac{(4\pi)^2}{130} \frac{M_{\text{GUT}}}{M_{\text{Pl}} c} \right]}$$

$\sigma(10^{\pm 7})$  FOR  $c=\pm 1$

$M_T$  VERY DEPENDENT ON  $c$

$\frac{\mu_3}{\mu_8}$  AND  $c$  CONSTRAINED

## ② X, Y GAUGE BOSONS

$$m_{\tilde{f}} \sim \frac{F}{M_{\text{GUT}}} \cdot \frac{\alpha_{\text{GUT}}}{4\pi}$$

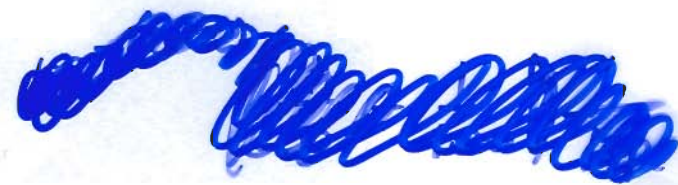
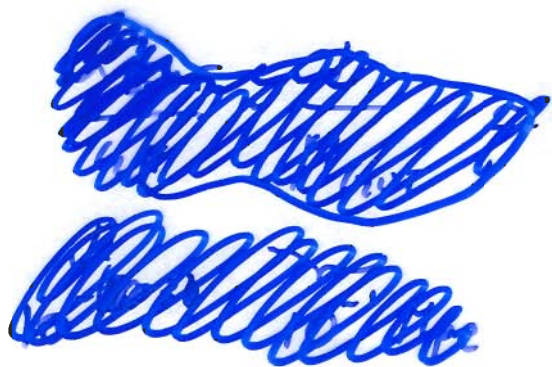
BUT IN OUR CASE

$$\frac{\alpha_{\text{GUT}}}{4\pi} \cdot \frac{1}{M_{\text{GUT}}} \ll \frac{1}{M_{\text{Pl}}}$$

X, Y GAUGE  
MEDIATED

GRAVITY  
MEDIATED

NEGLECTIBLE



$$\frac{\alpha_{\text{GUT}}}{4\pi} \cdot \frac{1}{M_{\text{GUT}}} \approx \frac{1}{10} \frac{1}{M_{\text{Pl}}}$$

↓                      ↓

$$\lesssim 10^{-2} \quad 10/M_{\text{Pl}}$$

(3)  $\sigma_3, \sigma_8$  FROM 244

$$u_{\tilde{f}} \approx \frac{\alpha}{4\pi} \frac{F}{u_i(\sigma)} \frac{\partial u_i(\sigma)}{\partial \sigma}$$

$i=3,8$

IN PRINCIPLE POSSIBLE:

$u_i(\sigma)$  POLYNOMIAL IN  $\sigma$

→ MAKE  $u_i(\sigma)$  SMALL

KEEPING  $\frac{\partial u_i(\sigma)}{\partial \sigma}$  LARGE

EXAMPLE:  $u(\sigma) = a\sigma^2 + b\sigma$

$$\Rightarrow a\langle\sigma\rangle^2 + b\langle\sigma\rangle \rightarrow 0$$

~~→~~ ↑ FINE-TUNE  $a, b$

$$\Rightarrow 2a\langle\sigma\rangle + b \neq 0$$

LOWER LIMIT FROM  $M_{\text{out}} < M_{\text{pe}}$

TYPICALLY SUBDOMINANT, MAYBE  
BORDERLINE

COMMENT:

WE ASSUMED LOW-ENERGY SUSY

$$\Lambda_{\text{MSSM}} \approx m_z$$

IF NOT

$$M_{\text{GUT}} = M_{\text{GUT}}^0 (\approx 10^{16} \text{ GeV})$$

$$m_T = m_T^0 (\approx 10^{16} \text{ GeV})$$

EVEN WITH

$$\frac{m_3}{M_{\text{GUT}}} = \left( \frac{m_z}{\Lambda_{\text{MSSM}}} \right)^{5/6}$$

$$m_3 \downarrow \quad \Lambda_{\text{MSSM}} \uparrow$$

$$\frac{m_g}{M_{\text{GUT}}} = \left( \frac{m_z}{\Lambda_{\text{MSSM}}} \right)^{1/2}$$

$$m_g \downarrow \quad \Lambda_{\text{MSSM}} \uparrow$$

(BUT LESS)

(GORAN SENJANOVIĆ  
TALK )

$$\Lambda_{\text{MSSM}} \rightarrow M_{\text{GUT}}$$

$$m_3 \rightarrow m_z$$

$$m_g \rightarrow 10^8 \text{ GeV (HALF WAY)}$$

# ④ HIGGS DOUBLETS

$$\int d^2\theta w_H(\sigma) \bar{H}H + h.c. \quad \leftarrow \text{FROM } \bar{H}(m + \lambda \Sigma + \frac{\Sigma^2}{M_{pe}} + \dots)H$$

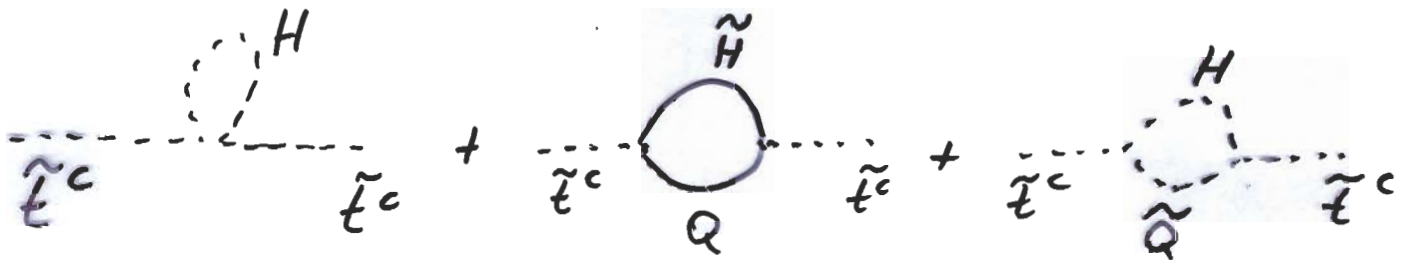
$$w_H(\langle \sigma \rangle) = \mu$$

FINE-TUNING  
(DT SPLITTING)

$$F \frac{\partial w_H(\langle \sigma \rangle)}{\partial \sigma} = B$$

EXTRA FINE-TUNING  
( $B \leq \mu^2$ )

1-LOOP



$$w_{\tilde{t}^c}^2 \simeq - \frac{y_t^2}{(4\pi)^2} B \quad \text{NEGATIVE!}$$

$H, \bar{H}$  CONTRIBUTION CANNOT DOMINATE  
(BUT NOT NECESSARILY NEGLIGIBLE)

# MESSAGE OF THIS TALK

- (1) TAKE YOUR FAVOURITE THEORY BEYOND THE STANDARD MODEL
  - MINIMAL
  - REALISTIC (  $G_{321}$  at low  $E$ ,  
correct spectrum,  
 $m_{\tilde{g}}$ ,  $V_{CKM}$ ,  $V_{PMNS}$  )
- (2) BREAK SUSY IN THIS MODEL (IF POSSIBLE)
- (3) FIND OUT THE MEDIATION
- (4) ANY PREDICTION ?  
(  $m_{3,8}$  LIGHT )

# ANOTHER EXERCISE

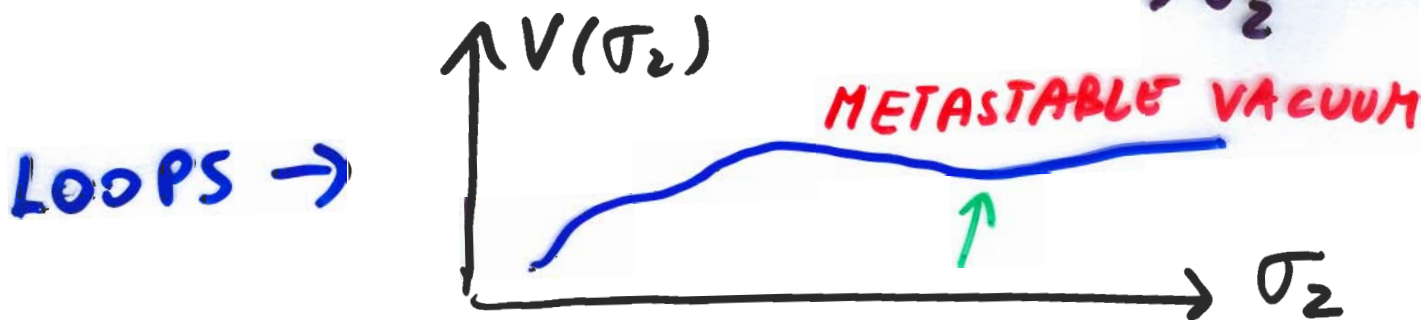
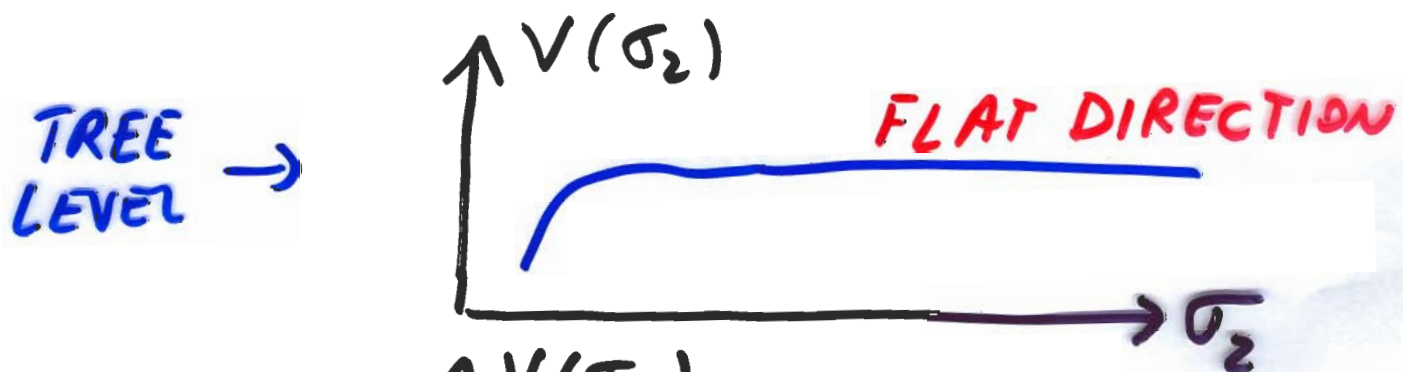
(1) RENORMALIZABLE  $SU(5)$   
(NO GRAVITY!)

(2)  $W = \xi X$  ;  $F_X = \xi$   
IN  $SU(5)$   $\Sigma_i \dots 24_i$

$$W = \text{Tr}(\Sigma_1^2 \Sigma_2)$$

$$\langle \Sigma_1 \rangle = v_1 \begin{pmatrix} 2 & & & & \\ & 2 & & & \\ & & -3 & & \\ & & & -3 & \\ & & & & -3 \end{pmatrix}$$

$$\Rightarrow W \simeq v_1^2 \sigma_2 \quad ; \quad F_2 \simeq v_1^2$$



# USE OF METASTABLE VACUA FOR SUSY BREAKING KNOWN LONG AGO

DIMOPoulos, DVALI, GIUDICE, RATTAZZI, 97

- REALISTIC EXAMPLES
- DECAY RATES

REVIVAL:

INTRILIGATOR, SHIH, SEIBERG, 06



SU(5) WITH TWO 24  
AND RENORMALIZABLE W

CAN SPONTANEOUSLY  
BREAK SUSY

- PROBLEMS

(1)  $3_w + 8_c$  LIGHT

GAUGE COUPLING UNIFICATION?

(2) DIFFICULT TO SPLIT  $D\bar{1}$

- EITHER FAST PROTON DECAY

- OR ~~DESTABILIZE~~ DESTABILIZE

SU(2) BREAKING

# CONCLUSIONS

① POSSIBLE TO USE MINIMAL  
(NON RENORMALIZABLE)  $SU(5)$  TO  
BREAK BOTH GUT AND SUSY

② PRICE TO PAY: FINE-TUNINGS

$$\rightarrow m_{\sigma} \ll M_{\text{GUT}} \quad (\text{NEW})$$

$$\rightarrow \mu \ll M_{\text{GUT}}$$

$$\rightarrow B \ll M_{\text{GUT}}^2$$

③ PREDICTIONS:

$\rightarrow M_{\text{GUT}}$  LARGE ( $d=6$  p decay  
NEGLECTIBLE)

( $d=5$  ALSO EASILY SAVED)

$\rightarrow$  GRAVITY MEDIATION DOMINATES